Influence of the neutron star 1E 161348–5055 in RCW 103 on the surrounding medium

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Abstract

We have carried out a study of the neutral hydrogen in the direction of the X-ray source 1E 161348–5055, a compact central object (CCO) located in the interior of the supernova remnant (SNR) RCW 103. The H I 21 cm line observations were carried out using the Australia Telescope Compact Array, complemented with single dish data from the Parkes radio telescope to recover information at all spatial scales. We derive a distance to RCW 103 of 3.3 kpc, in agreement with previous distance measurements. We have also detected a small hole in the H I emission which is positionally and kinematically coincident with the location of the CCO which confirms the association between the SNR and the CCO. This is the third case of a depression in H I emission seemingly associated with CCOs in SNRs. The characteristic parameters of the holes such as their size, eccentricity and evacuated mass are similar in all three cases. We estimate the absorbing H I column density towards 1E 161348–5055 to be $\sim 6 \times 10^{21}$ cm^{-2}, a value compatible with a blackbody solution for the CCO X-ray emission. However, the implied brightness temperature is very high compared to most neutron stars. Moreover, the strong long-term variability in X-rays favours the hypothesis that 1E 161348–5055 is an accreting binary source rather than an isolated, cooling neutron star. An analysis of the continuum image obtained at 1.4 GHz from these observations shows no trace of a pulsar wind nebula around 1E 161348–5055, in spite of it being a young object.

Keywords: stars: neutron – supernova remnants – ISM: individual: RCW 103 – X–rays: individual: 1E 161348–5055 – spectral lines: neutral hydrogen.

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1 Introduction

X-ray observations from the past few years have revealed the existence of a large variety of Galactic point-like sources, many of them identified with neutron stars but presenting very different observational properties. These unresolved sources, which appear either isolated or in the interior of supernova remnants (SNR), have no radio counterpart and very high X-ray to optical flux ratios. About half of these sources show X-ray pulsations with periods between 6 and 12 seconds (anomalous X-ray pulsars, AXP) and can even present sporadic strong γ-ray emission (soft gamma-ray repeaters, SGR). The rest of these sources are generally called either isolated neutron stars (INS) if they are not associated with any other object in the sky, or central compact objects (CCO) if they appear projected on the interior of a SNR (see a recent review by Pavlov et al. 2002).

The nature of most CCOs is still unclear. Brazier & Johnston (1999) explain them as normal pulsars with unfavourable radio beaming, while the soft thermal X-rays would be easily detected due to its almost isotropic emission. On the other hand, Vasisht et al. (1997) and Frail (1998) conclude that CCOs are neutron stars (NSs) born with long initial periods and high magnetic fields ($B > 10^{14}$ G), and thus would be related to AXPs and SGRs. However, Geppert, Page & Zannias (1999) suggest that CCOs are fast-spinning, weakly-magnetized sources. The X-ray emission from CCOs is generally explained as thermal radiation from cooling NSs (e.g Zavlin, Trümper & Pavlov 1999), with typical temperatures of a few 10^6 K, as inferred from their thermal-like spectra.

In a recent H I study towards the bilateral SNR G296.5+10.0, Giacani et al. (2000) found that the associated CCO, 1E 1207.4–5209, lies near the center of a small H I depression located at the same systemic velocity as the SNR. The authors propose that the depression is due to self-absorption of a cool H I cloud just foreground to a hotter volume of gas surrounding the CCO, and heated by its X-ray flux. We have begun a systematic search for similar traces in the neutral gas around other CCOs. The observations towards the X-ray point source RX J0822–4300 in Puppis A revealed an H I structure consisting of a nearly circular minimum coincident with the CCO plus two aligned lobe-like depressions that appear to emerge from the CCO (Reynoso et al. 2003). The two lobes appear to have been formed by a combination of the proper motion of the CCO and the ejection of a collimated outflow. In this paper, we present the results obtained for 1E 161348–5055, the CCO associated with the SNR RCW 103 (G332.4–0.4).

At radio wavelengths, RCW 103 appears as an almost complete, circular $8'$ diameter shell (Caswell et al. 1980). High-resolution radio polarimetric data reveal an approximately east-west alignment of the magnetic field which could be helping to limit the expansion north-south (Dickel et al. 1996). The brightening of the rim in the northern and southern sides, that hints for an incipient bilateral barrel shape for this SNR, could be related to this frozen-in magnetic field. Based on H I absorption measurements at 21 cm, Caswell et al. (1975) suggest a distance of 3.3 kpc. Optical filaments are seen toward the brighter regions of the radio shell (van den Bergh, Marscher & Terzian 1973;

Ruiz 1983), and observations at infrared wavelengths show evidence for interaction with a dense interstellar medium (ISM), particularly to the south (Oliva, Moorwood & Danziger 1990; Burton & Spyromilio 1993; Oliva et al. 1999). An optical expansion study of RCW 103 (Carter, Dickel & Bomans 1997) indicates that this SNR is probably about 2,000 years old, based on an assumed distance of 3.3 kpc. However, optical extinction studies suggest distances around 6.5 kpc (Leibowitz & Danziger 1983; Ruiz 1983).

Soft X-ray emission from RCW 103 was first detected with the Einstein Observatory (Tuohy et al. 1979), indicating a close correlation with the non-thermal radio emission. Also, a faint, point X-ray source, 1E 161348–5055 , was located near the centre of the SNR (Tuohy & Garmire 1980). The lack of optical or radio counterparts led Tuohy et al. (1983) to propose that this point source was a thermally radiating CCO. However, subsequent observations with the Einstein IPC and ROSAT failed to confirm the existence of this source (Becker et al. 1993). Finally, Gotthelf, Petre & Hwang (1997) detected hard X-rays from the elusive CCO using ASCA. They found that the spectral characteristics were incompatible with a simple cooling NS model. Follow-up observations (Gotthelf, Petre & Vasisht 1999) confirmed that 1E 161348–5055 manifests long-term variability, explaining the non-detections after its discovery. There are strong indications that 1E 161348–5055 is part of a binary system (see §4.1).

In this paper, we present radio contiuum at 1.38 GHz and H I λ21 cm observations carried out with the Australia Telescope Compact Array (ATCA) towards the SNR RCW 103.

2 Observations and data reduction

Interferometric observations were obtained with the ATCA during one session of 12 h with the 750A array (baselines from 76.5 to 735 m), on 2002 January 22, and one session of 12 h with the EW 367 array (baselines from 46 to 367 m) on 2002 April 1. The antennas were pointed at RA= $16^{\mathrm{h}}17^{\mathrm{m}}30\overset{\mathrm{s}}{.}0$, Dec.= $-51°0'0''$ (J2000). A correlator configuration of 1024 channels, covering a total bandwith of 4 MHz centered at 1420 MHz, was used. The corresponding velocity resolution at this frequency is 1 km s^{-1}. Continuum data were obtained simultaneously with a bandwidth of 128 MHz centered at 1384 MHz. The source PKS B1934–638 was used for primary flux density calibration, while phases were calibrated with PKS B1657–56.

The data were processed with the MIRIAD software package (Sault, Tauben & Wright 1995). To subtract the continuum component from the H I data set, a linear baseline was fitted to 370 line-free channels. The final H I cube was constructed with a $15''$ cell size, covering an area of $64' \times 64'$ and keeping 350 channels from –150 to +138 km s^{-1}. Sidelobes were supressed over an area of $34' \times 34'$. The image was cleaned and convolved with a $50''$ Gaussian beam. The continuum image was constructed with the same geometry and angular resolution as the H I data. After cleaning and restoring with a $50''$ Gaussian beam, a sensitivity of $\sigma \simeq 5.5$ mJy beam^{-1} was achieved.

To recover structures at shorter spatial frequencies, the ATCA H I data were combined in the *u,v* plane with single dish data from the Southern Galactic Plane Survey

Figure 1: Radio continuum image centred at RCW 103, obtained with the ATCA at 1384 MHz. The gray scale is shown on top of the image in units of Jy beam^{-1}. The beam, $50'' \times 50''$, is plotted in the bottom right corner. The noise level is 5.5 mJy beam^{-1}.

Figure 2: (a) Average H I emission profile towards RCW 103 using interferometric and single dish data. (b) Absorption profile towards RCW 103 based on interferometric ATCA data. (c) Same as (b) but excluding baselines shorter than 1 kλ. In all cases, the velocity resolution is 1 km s^{-1}.

(SGPS; McClure–Griffiths et al. 2001) obtained with the Parkes telescope. In order to match the intensity units in both sets, a factor of 257.5 was applied to the ATCA data cube to convert flux densities to brightness temperatures. No tapering was applied to the low resolution cube. The rms of the combined image is 2 K per channel.

3 Results

3.1 Absorption study

To constrain the systemic velocity of RCW 103, we analyzed the H I absorption against the strong radio continuum emission from the remnant. The radio contimuum image obtained with ATCA at the frequency of 1384 MHz in a field of about $45'$ around RCW 103 is shown in Figure 1. According to the equation of radiative transfer, the emission T_{b_v} at a velocity v measured by an interferometer is given by

$$T_{b_v} = T_s(1 - e^{-\tau_v}) + T_c e^{-\tau_v}, \tag{1}$$

where T_s and τ_v are the spin temperature and the optical depth of the intervening H I gas, and T_c is the background continuum emission. Since we have subtracted the radio continuum from out H I data, then eq.(1) can be reduced to

$$T_{L_v} = (T_s - T_c)(1 - e^{-\tau_v}), \tag{2}$$

where $T_{L_v} = T_{b_v} - T_c$. Equation (1) can be simplified to a single unknown if $T_c \gg T_s$. This condition can be accomplished by using only interferometric data, since in such way the extended H I emission is filtered out.

Figure 2a shows an average H I emission profile towards RCW 103 combining interferometric plus single dish data, while Fig. 2b shows an average profile of $e^{-\tau_v}$ over the same area, computed as described above (eq. (2)). Strong absorption peaks appear at –43, –17, +1, and +34 km s^{-1}. All but the last feature were detected by Caswell et al. (1975) based on data obtained with the Parkes interferometer with a 2 km s^{-1} spectral resolution. We do not believe the absorption at +34 km s^{-1} to be real because (a) absorption features should appear also from –43 km s^{-1} to the tangent point at –115 km s^{-1}, and (b) the distance to the SNR would be $\sim$ 20 kpc, which gives an unrealistic size and expansion velocity.

Figure 3: Gray-scale and contour image of the average H I emission between –46.1 and –53.5 km s^{-1} towards RCW 103. The rms is 2K. The beam, $50'' \times 50''$, is plotted as a white open circle in the bottom left corner. The brightness-temperature scale is indicated on top of the image in units of K, while the contour levels vary from 68 K in steps of 6 K. For comparison, a few representative contours of the radio continuum emission are included as white lines. The cross indicates the position of 1E 161348–5055 as given by Garmire et al. (2000a).

For comparison, we obtained ATCA spectra towards the strong Galactic radio sources present in the field G332.7–0.6 and G332.2–0.4 (Fig. 1), placed within $15'$ of RCW 103, and found that absorption features appear only in the velocity ranges from 0 to –50 and to –58 km s^{-1} respectively.

To solve this puzzle, we followed the method employed by Dickey et al. (2003). We constructed a new H I cube without subtracting the radio continuum and removing all baselines shorter than 1 kλ (210 m). In this way, all H I structures larger than $\sim 4'$ are filtered out. The term involving T_s in eq. (1) can thus be neglected. The data cube was further weighted by the radio continuum. We then computed $e^{-\tau_v}$ using eq. (1). The resulting profile is plotted in Figure 2c.

This filtering method succeeded in removing most of the H I emission, as is apparent by comparing Figs. 2b and 2c. The absorption feature at –43 km s^{-1} is confirmed, while the feature at +34 km s^{-1}has disappeared. This implies that the latter was produced by self absorption on a scale between $4'$ and $\sim 30'$.

The absorption feature at –43 km s^{-1} corresponds to a lower distance limit of 3.1 kpc according to the Galactic rotation model of Fich, Blitz & Stark (1989). We do not see any absorption against the emission peak at ~ -75 km s^{-1}, even though it has a brightness temperature of ~ 80 K. Thus the upper distance limit is 4.6 kpc. The lower distance coincides with the near side of the Scutum-Crux arm (Georgelin & Georgelin 1976), and will be assumed throughout the paper.

3.2 H I associated with RCW 103

We estimated the H I column density towards 1E 161348–5055 by integrating the foreground H I brightness temperature of the interferometric plus single-dish data up to –43 km s^{-1}. The value obtained depends on the lower limit adopted for the velocity interval in the integration since, although the local gas is supposed to lie at a systemic velocity of 0 km s^{-1}, turbulence may cause departures of about 7 km s^{-1} from this value. Besides, due to the distance ambiguity, it is possible that some background gas is included in the integration, leading to an overestimation of $N_{\rm H}$. At $v \simeq -10$ km s^{-1}, the line of sight crosses the far side of the Scutum-Crux arm. Taking into account these considerations, we integrate the brightness temperature between $\sim +3.5$ and -43 km s^{-1} and estimate the H I column density to be $N_{\rm H} \lesssim 6 \times 10^{21}$ cm^{-2}. The implicances of this determination will be discussed in §4.

An inspection of the H I images at velocities around –43 km s^{-1} revealed that 1E

161348–5055 lies inside a local H I depression which is present in all channels between –46.1 and –53.5 km s^{-1}. In Figure 3, an image of the average H I emission within this velocity interval is shown. The H I depression, which attains its minimum at RA= $16^{\rm h}17^{\rm m}34\overset{\rm s}{.}8$, Dec.= $-51°2'40''$ (J2000), is elongated, with a minor-to-major axis ratio of ~ 0.7, and has a mean diameter of $\sim 64''$ (1 pc at a distance of 3.1 kpc). 1E 161348–5055 is $20''$ (0.3 pc) away from the centre of the hole. We note that due to the proximity to the Galactic Plane, our data show several similar H I depressions at different locations and velocities. However, the coincidence in position and velocity makes an association between the CCO and this H I feature very likely.

In what follows, we analyze possible origins for this H I void. If the H I minimum is produced because of a real absence of H I, then the missing mass is estimated to be 0.3 $M_\odot$. If instead it is produced by hot H I gas self-absorbed by a cooler foreground, as proposed for the CCO 1E 1207.4–5209 (Giacani et al. 2000), then we can follow the method described by Schwarz et al. (1995) to estimate an upper limit for the temperature of the hot neutral hydrogen gas.

If self-absorption is considered, then eq. (1) can be re-written as (cf. eq. (3) in Schwarz et al. 1995)

$$T_{L_v} = (T_s - T_{bg})A_v - T_c A_v, \tag{3}$$

with $A_v = (1 - e^{-\tau_v})$, and where T_{bg} is the temperature of the background hot H I. Assuming that A_v is uniform across the continuum source, then T_{L_v} is a linear function of T_c with slope $(-A_v)$ and a zero offset $(T_s - T_{bg})A_v$. The H I hole attains its minimum emission at $v = -46.1$ km s^{-1}. We compared the brightness temperature of the ATCA H I data at this velocity with the continuum emission and fitted a straight line to the distribution. To avoid confusion possibly introduced by regions of low continuum emission, all values of T_c under 350 mJy beam^{-1} were clipped. We obtained that $A_v \simeq 0.195$, which implies that $\tau_v \simeq 0.2$. To estimate the spin temperature, it must be noticed that for a single dish measurement towards a region with no continuum emission, the brightness temperature of the line is $T_{L_v} = T_s A_v$ (Schwarz et al. 1995). We averaged the Parkes data in a box around RCW 103 (excluding the contribution from the SNR) and obtained that $<T_{L_v}> \leq 100$ K, hence $T_s \leq 20$ K. Finally, at the location of 1E 161348–5055, the interferometric data give $T_c = 44$ K and $T_{L_v} = -30$ K. Thus, applying eq. (3) to the emission towards the CCO, the hot H I gas in the hole is found to have a temperature of ≤ 130 K. Such a temperature is unrealistically low at the interior of a SNR, therefore we discard the possibility that the H I minima observed around CCOs are produced by self-absorption.

3.3 Radio continuum emission

The rotational energy of pulsars is dissipated via a magnetized relativistic wind composed of electrons and positrons. The shock front between this wind and the ambient medium can give rise to a synchrotron emitting bubble (known as a pulsar wind nebula, PWN). The detectability of such a PWN depends on the density of the ambient medium and the pulsar parameters. The majority of pulsars do not have detectable PWN (Gaensler et al. 2000). For a very young pulsar inside an SNR, the luminosity

Table 1: Parameters of the H I depressions around CCOs

CCO (associated SNR)	1E 1207.4–5209[a] (G296.5+10.0)	RX J0822–4300 [b] (Puppis A)	1E 161348–5055 [c] (RCW 103)
mean angular diameter (′)	5.3	2.3	1.1
mean linear diameter (pc)	3.2	1.5	1.0
minor/major axis ratio	0.8	0.8	0.7
CCO offset from centre (″)	30	37	20
CCO offset from centre (pc)	0.3	0.4	0.3
missing mass ($M_\odot$)	–	0.1	0.3

[a]Giacani et al. (2000)
[b]Reynoso et al. (2003)
[c]Present work.

Figure 4: Gray-scale and contour image of RCW 103 at the radio continuum frequency of 1384 MHz. The gray scale is shown on top of the image in units of Jy beam^{-1}. The contours are plotted in steps of 7.5% of the peak intensity, 768 mJy beam^{-1}, starting at 15%. For clarity purposes, white lines are used over dark background. The cross shows the position of 1E 161348–5055 as given by Garmire et al. (2000a). The noise level is 5.5 mJy beam^{-1}.

L of the PWN is a strong function of the initial period of the pulsar, P_0 ($L \propto P_0^{-5}$; Reynolds and Chevalier 1984).

We do not detect a PWN at the position of 1E 161348–5055. Rather, the radio continuum emission in RCW 103 has a minimum near the CCO (Figure 4). To investigate if there is a PWN around the CCO that could be beam-diluted at the resolution of 50″, we constructed another continuum image using only the longest baselines. This image, with a resolution of $6\farcs2 \times 4\farcs6$, shows no emission at the position of the CCO down to a level of $5\sigma = 1$ mJy beam^{-1}. This limit is not very constraining - it may imply either that the spin period of 1E 161348–5055 is not particularly short or that the PWN is not well confined in the interior of the SNR.

4 Discussion

4.1 Spectral model

Gotthelf et al. (1997) investigated different fits to the X-ray spectrum of 1E 161348–5055 using a nonequilibrium ionization plasma model combined alternatively with a blackbody, a power-law, and a thermal bremsstrahlung continuum component. The best fits were achieved with H I column densities of 5, 31, and 16 $\times 10^{21}$ cm^{-2} respectively, although a fixed value of $N_H = 6.7 \times 10^{21}$ cm^{-2} also gives reasonable fits in all three cases. The value of N_H that we estimated from the present observations, $N_H \leq 6 \times 10^{21}$ cm^{-2}, clearly favours the blackbody solution in this case.

However, since 1997, there is growing evidence to indicate that 1E 161348–5055 might be an accreting binary system. Chandra and ASCA observations revealed a sinusoidal light curve with a period of ~ 6.4 hr (Garmire et al. 2000a). Even though subsequent observations failed to detect the modulation again (Garmire et al. 2000b) due to pile-up in the detectors from the very high X-ray flux, the detection of two partial dips separated by 180° in phase in the light curve, and a possible near-IR counterpart about $2''$ away from the nominal Chandra position, suggest that 1E 161348–5055 may be powered by accretion likely from a low-mass companion in a binary system (Sanwal et al. 2002). Finally, Becker & Aschenbach (2002) discovered an eclipse 3 h after the start of their observations using XMM-Newton data. This would make 1E 161348–5055 the first accreting binary detected inside a SNR. Becker & Aschenbach (2002) also obtained an X-ray spectrum towards 1E 161348–5055 and found that a simple blackbody model describes the data almost up to ~ 5 keV but not beyond. They obtain a best fit with a double blackbody, but the required $N_{\rm H}$ of $\leq 18 \times 10^{21}$ cm^{-2} is incompatible with our results.

We note that all the blackbody models imply temperatures as high as $8-10\times10^6$ K (Becker & Aschenbach 2002) and a luminosity of $\sim 10^{34}$ ergs s^{-1} (Gotthelf et al. 1997), at least twice the values predicted by theoretical models (see Gotthelf et al. 1997 and references therein). In general, fitting NS models to the compact sources inside SNRs presents the problem that the derived parameters (temperature and luminosity) are far larger than normally expected (Lloyd, Hernquist & Heyl 2002). In addition, estimates of X-rays emitting areas based on blackbody models are usually too small to be the size of a NS (e.g. Becker & Aschenbach 2002). One possible solution is proposed by Heyl & Hernquist (1998), who show that a hot blackbody can be mimicked if the cooling NS model includes an ultramagnetized star (B $\sim 10^{15}$ G) with an accreted hydrogen atmosphere. This explanation was also invoked for other CCOs, like RX J0822–43 in Puppis A (Zavlin et al. 1999) and 1E 1207.4–5209 in G296.5+10.0 (Zavlin et al. 1998). Even in non-magnetic atmospheres, the effect of H or He atmospheres is to deviate the high energy spectrum in such a way that the effective temperatures as given by blackbody models could be overestimated by a factor of 2 (Zavlin, Pavlov & Shibanov 1996).

In summary, it seems most likely that 1E 161348–5055 is a binary system and that fits to the X-ray spectrum need to be re-visited. On the other hand, it is still possible that this system is a single neutron star. The constraints we have put on $N_{\rm H}$ show that it is very hot for its age, and may indicate exotic atmospheric processes at work.

4.2 The H I depression

The present study reveals the third case in which a CCO associated with a SNR is located inside an H I depression. The other two cases are 1E 1207.4–5209 in the SNR G296.5+10.0, and RX J0822–4300 in Puppis A. In Table 1 we compare the main parameters derived for the three H I depressions. In G296.5+10.0, only the size and velocity width are given in Giacani et al. (2000), thus the remaining parameters, where possible, have been estimated directly on the images presented in their paper.

In computing the minor/major axis ratio, the beam elongation was taken into account.

A number of similarities can be found between the three cases. All H I cavities appear to have the same elongation, and the CCOs are off-centred by similar distances. In the two cases where the missing mass was computed, a similar value was obtained. On the other hand, the size of the H I hole in G296.5+10.0 is larger than the other two holes by more than a factor of 2. This discrepancy may be related to the beam size of the observations, which is almost 4 times larger than those of the Puppis A and RCW 103 data. The main difference between the three cases is given by Puppis A, which presents two lobe-like H I depressions emerging from its associated CCO at the systemic velocity of the SNR (Reynoso et al. 2003). The present observations do not reveal any similar morphology around 1E 161348–5055.

In what follows we will analyze if the H I depression around 1E 161348–5055 can be a swept up hole. The rate of energy needed to set the missing mass in the cavity into motion with velocity $V = r/t$ is

$$\dot{E}_k = 3 \times 10^{47} M r^2 t^{-3} \text{ erg s}^{-1}, \tag{4}$$

where M is the evacuated mass in $M_\odot$, r is the radius of the hole in pc, and t is the age of the CCO, equal to the age of the host SNR, in years. Replacing in eq. (4) the radius and missing mass computed in §3, and assuming that the age of RCW 103 is 2,000 yrs (Carter et al. 1997), we obtain $\dot{E}_k = 3^{+20}_{-2} \times 10^{36}$ erg s^{-1}, where the quoted errors allow for ages of 1,000 and 3,000 yrs old. The spin down energy loss has not been measured for 1E 161348–5055, but if it is similar to the values observed in other CCOs (energy loss rates between $\sim 1 \times 10^{36}$ erg and 1.5×10^{37} erg; Slane et al. 1997, Brazier & Johnston 1999), then a complete conversion of the rotational energy into kinetic energy of the surrounding medium can account for the observed H I hole.

An alternative possibility is that the depression does not contain swept up, low density gas, but is filled with H I gas heated up by the CCO at temperatures higher than the surroundings. In that case, the depression should contain enhanced ionized hydrogen (i.e. it would form a small H II region), and this is not observed in the radiocontinuum image. We would, however, expect to see infra-red emission from such a region but a search in mid- and near-infrared wavelengths using data from the Midcourse Space Experiment (MSX) and the Two Micron All Sky Survey (2MASS), yielded negative results. In conclusion, we find that the swept-up cavity provides a more convincing explanation for the observed H I minimum.

5 Conclusions

In this paper, we present the third case in which a CCO lies at a local H I minimum at a velocity compatible with the systemic velocity of the host SNR. We have shown that self-absorption does not provide a satisfactory explanation for this kind of features, as was proposed for 1E 1207.4–5209 (Giacani et al. 2000). Instead, it is possible that the H I depression is a swept up hole, where $\sim 0.3 M_\odot$ of neutral gas has been evacuated. We have found a number of similarities between the three H I holes around CCOs

detected so far, like the elongation, the missing mass, and the off-centred position of the CCOs. We did not detect any synchrotron nebula around the CCO down to a level of 1 mJy beam^{-1}.

The present data allow us to constrain the H I column density to be $N_H \sim 6 \times 10^{21}$ cm^{-2}. This column density favours the blackbody model of Gotthelf et al. (1999), however the derived temperature is too high to be explained by standard NS cooling models. Instead, it appears most likely that 1E 161348–5055 is an accreting binary, and may constitute the first case in which such a system occurs in the interior of a SNR.

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